

# Trace-distance measure of Coherence

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# Joint works...

Collaborators: Maciej Lewenstein, Swapan Rana, and Andreas Winter

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## PHYSICAL REVIEW A

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### Trace-distance measure of coherence

Swapan Rana, Preeti Parashar, and Maciej Lewenstein  
Phys. Rev. A **93**, 012110 – Published 12 January 2016



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### Logarithmic coherence: Operational interpretation of $\ell_1$ -norm coherence

Swapan Rana, Preeti Parashar, Andreas Winter, and Maciej Lewenstein  
Phys. Rev. A **96**, 052336 – Published 27 November 2017



# Outline

- ▶ **Introduction to coherence measures**
  - Model: Incoherent Operations (IO)
  - Coherence measures: basic requirements
- ▶ **Distance-based measures**
  - $\ell_p$  - and Shatten- $p$ -norm based measures?
  - Neither is strong monotone for  $p > 1$
- ▶ **Trace distance measure of coherence**
  - Strong monotone for Qubits and  $X$ -states
  - But not for all states
  - What about entanglement?
- ▶ **Operational interpretations of  $C_{\ell_1}$** 
  - Equals to Negativity of the MC state
  - Optimum  $C_r$  for a given  $C_{\ell_1}$ ?
- ▶ **Conclusion**

# Coherence Theory

## Resource Theory of Coherence [Baumgratz *et al.*, PRL (2014)]

### ► Free states: Diagonal in chosen basis

- 'Preferred', or 'Chosen' basis  $\{|i\rangle\}$  of the corresponding Hilbert space  $\mathcal{H} = \mathbb{C}^d$ .
- Set of 'incoherent', or 'free' states:

$$\mathcal{I} = \left\{ \delta \mid \delta = \sum_i p_i |i\rangle\langle i| \right\}.$$

### ► Free Operations: Incoherent Operations (IO)

- An operator  $K$  is 'incoherent' if  $K\delta K^\dagger \in \mathcal{I}$ ,  $\forall \delta \in \mathcal{I}$ .  
 $\implies K$  can have at most one non-zero entry in any column.
- A CPTP map  $\Lambda$  is 'free' or 'incoherent' if  $\exists$  incoherent Kraus operators  $\{K_n\}$  for  $\Lambda$ .

# Coherence measures

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### Coherence Monotones

Any function  $C: \mathcal{B}(\mathcal{H}^d) \rightarrow \mathbb{R}^+$  is a coherence measure if it satisfies

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- 3a. Monotonicity:  $C(\Lambda_{\text{ICPTP}}[\rho]) \leq C(\rho)$ ,  $\forall \rho$  and incoherent  $\Lambda$ .

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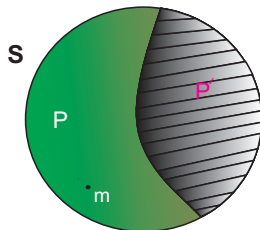
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- 3a. Monotonicity:  $C(\Lambda_{\text{ICPTP}}[\rho]) \leq C(\rho)$ ,  $\forall \rho$  and incoherent  $\Lambda$ .
- 3b. Strong Monotonicity:  $\sum p_n C(\rho_n) \leq C(\rho)$ ,  $\rho_n = K_n \rho K_n^\dagger / p_n$ ,  $p_n = \text{Tr}[K_n \rho K_n^\dagger]$ , for all incoherent  $\{K_n\}$ . Allows selective measurements or stochastic operations. Harder to verify:  
3b. + 2.  $\Rightarrow$  3a.

## Distance-based measures

### How to quantify (measure)?

Let  $m \in S$  and  $P$  be a *property* regarding the members of  $S$ . The usual way to quantify how much of the property  $P$  is contained in  $m$  is to determine the distance of  $m$  from the set of all  $x \in S$  which does not have the property  $P$ .



## Entanglement from contractive metric [Vedral and Plenio, PRA 1998]

If a *distance function*  $\mathcal{D}$  satisfies

- i. Positivity:  $\mathcal{D}(\rho, \sigma) \geq 0 \quad \forall \rho, \sigma$ , with equality iff  $\rho = \sigma$ ,
- ii. Contractivity:  $\mathcal{D}(\Lambda[\rho], \Lambda[\sigma]) \leq \mathcal{D}(\rho, \sigma)$  for all CPTP map  $\Lambda$ ,

then an entanglement measure can be defined through this distance as

$$E(\rho) = \inf_{\sigma \in \{\text{Separable states}\}} \mathcal{D}(\rho, \sigma).$$

✓ REE :

$$\mathcal{D}(x, y) = S(x\|y) := \begin{cases} \text{Tr}(x \log x - x \log y), & \text{if support } x \subseteq \text{support } y \\ +\infty, & \text{otherwise} \end{cases}$$

✓ Bures metric:  $\mathcal{D}(x, y) = 2 - 2\sqrt{F(x, y)}$ , where

$$F(x, y) := [\text{Tr}\{\sqrt{\sqrt{x}y\sqrt{x}}\}]^2$$

✓ Trace distance:  $\mathcal{D}(x, y) = \|x - y\|_1 := \text{Tr}|x - y|$

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✓ Trace distance:  $\mathcal{D}(x, y) = \|x - y\|_1 := \text{Tr}|x - y|$

✗ Hilbert-Schmidt distance:

$$\mathcal{D}(x, y) = \|x - y\|_2 \text{ is not-contractive [Ozawa, PLA 2000].}$$

## Measures from $\ell_p$ - and Schatten- $p$ -norm ( $C_{\ell_p}, C_p$ )

- For  $X \in \mathbb{C}^{m \times n}$ ,  $p \in [1, \infty)$ , the norms are defined by

$$\|X\|_{\ell_p} := \left( \sum_{i,j} |x_{ij}|^p \right)^{1/p},$$

$$\|X\|_p := (\text{Tr}|X|^p)^{1/p} = \left( \sum_i \sigma_i^p \right)^{1/p}, \quad \sigma = \lambda \left( \sqrt{X^\dagger X} \right).$$

- The induced distance functions  $C_{\ell_p}$  and  $C_p$  satisfies 1. and 2.
- So we have to check only 3: the (strong) monotonicity.

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- Both are monotone on qubits. Is it so in higher dim?

☕ For all  $p \in (1, \infty)$  neither  $C_{\ell_p}$  nor  $C_p$  is a monotone [PRA (2016)].



# None of $C_{\ell_p}, C_p$ is a monotone for $p > 1$

Counterexample from any  $\rho \notin \mathcal{I}$

Let  $K_i = |0\rangle\langle i-1|$ ,  $i = 1, 2, \dots, d$  and consider the incoherent channel  $\Lambda = \{\tilde{K}_i\}$  with the Kraus operators  $\tilde{K}_i = \mathbb{1} \otimes K_i$ . Then we have

$$\begin{aligned} C_p(\Lambda[\rho \otimes \mathbb{1}/d]) &= C_p(\rho \otimes |0\rangle\langle 0|) \\ &= C_p(\rho) \\ &> C_p(\rho \otimes \mathbb{1}/d). \end{aligned}$$

The inequality follows from

$$C_p(\rho \otimes \mathbb{1}/d) \leq \|\rho \otimes \mathbb{1}/d - \delta^* \otimes \mathbb{1}/d\|_p = C_p(\rho) \|\mathbb{1}/d\|_p < C_p(\rho).$$

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As  $C_{\ell_p}(\rho \otimes \mathbb{1}/d) = d^{1/p-1} C_{\ell_p}(\rho) < C_{\ell_p}(\rho)$ ,  $C_{\ell_p}$  also violates monotonicity.

Hence  $C_{\ell_p}, C_p$  is not strong monotone for  $p > 1$

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$$\rho = \frac{1}{4} \begin{pmatrix} 1 & 0 & a & 0 \\ 0 & 1 & 0 & b \\ \bar{a} & 0 & 1 & 0 \\ 0 & \bar{b} & 0 & 1 \end{pmatrix}, K_1 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}, K_2 = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix}.$$

Then

$$\rho_1 = \frac{1}{4p_1} \begin{pmatrix} 1 & a \\ \bar{a} & 1 \end{pmatrix} \oplus 0, \rho_2 = \frac{1}{4p_2} \begin{pmatrix} 1 & b \\ \bar{b} & 1 \end{pmatrix} \oplus 0, p_1 = p_2 = \frac{1}{2}.$$

The strong monotonicity (for both  $C_{\ell_p}$  and  $C_p$ ):  $(|a| + |b|)^p \leq |a|^p + |b|^p$ .

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$$(x + y)^{1+\epsilon} = (x + y)(x + y)^\epsilon = x(x + y)^\epsilon + y(x + y)^\epsilon > x \cdot x^\epsilon + y \cdot y^\epsilon = x^{1+\epsilon} + y^{1+\epsilon}$$

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Strong monotonicity does not hold for all  $ab \neq 0$  and  $p \in (1, \infty)$ .

The only remaining case  $p = 1$ :  $C_{\ell_1}$  and  $C_{\text{Tr}}$

$$C_{\ell_1}(\rho) := \min_{\delta \in \mathcal{S}} \|\rho - \delta\|_{\ell_1} \qquad C_{\text{Tr}}(\rho) := \min_{\delta \in \mathcal{S}} \|\rho - \delta\|_1$$

$$= \sum_{i \neq j} |\rho_{ij}| \qquad = ?$$

- $C_{\ell_1}$  is a (strong) monotone [Baumgratz *et al.*, PRL (2014)].
- $C_{\text{Tr}}$  is a monotone, as  $\|\cdot\|_1$  is contractive under any CPTP map.

Semidefinite Program (SDP) for  $C_{\text{Tr}}(\rho)$  [PRA (2016)]

$$\begin{aligned} & \text{Minimize} && \text{Tr}(P + N) \\ & \text{subject to} && \begin{cases} P - N = \rho - \delta, \\ \text{Tr} \delta = 1, \\ \delta \in \mathcal{S} \\ P, N, \delta \geq 0. \end{cases} \end{aligned}$$

For pure states, a better SDP: [Chen *et al.*, PRA (2016)].

## $C_{Tr}$ is monotone but not strong monotone

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Strong coherence monotone: any convex and monotone  $C$  is strong monotone iff  $C(p_1\rho_1 \oplus p_2\rho_2) = p_1C(\rho_1) + p_2C(\rho_2)$ .

☛ For  $|\Phi_d\rangle = \frac{1}{\sqrt{d}} \sum |i\rangle$ ,  $C_{\text{Tr}}(|\Phi_d\rangle) = 2(1 - 1/d)$ . Choosing  $\rho_1 = |\Phi_2\rangle$ ,  $\rho_2 = |\Phi_3\rangle$ ,  $p_1 = p_2 = 1/2$ , shows that  $C_{\text{Tr}}$  is not a strong monotone.

## What about $E_{\text{Tr}}$ ?

$$E_{\text{Tr}}(\rho) := \min_{\sigma \in \{\text{Separable states}\}} \|\rho - \sigma\|_1$$

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Same for pure states [Chen *et al.*, PRA (2016)]:

$$C_{\text{Tr}}\left(\sum \sqrt{\lambda_i} |i\rangle\right) = E_{\text{Tr}}\left(\sum \sqrt{\lambda_i} |ii\rangle\right)$$

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$E_{\text{Tr}}$  is not a strong monotone [Qiao *et al.*, arXiv: 1710.04447 (2017)]!

## Facts about $C_{\ell_1}$

- Captures the simple intuitive idea that superposition corresponds to off-diagonals.
- Easily computable even in the time-dependent case where the evolved density matrix can not be diagonalized, so eigenvalues and  $C_r$  become non-computable.
- Physical: Success probability of unambiguous state discrimination in interference experiments, 'which-path information' about a particle inside a multipath interferometer [Bagan *et al.* PRL (2016)].
- No conspicuous role in entanglement theory

## $C_{\ell_1}$ for pure states

- ☞ **Equals to Negativity:** For pure states,  $C_{\ell_1}$  is negativity of the corresponding bipartite state,

$$C_{\ell_1}(|\psi\rangle := \sum \sqrt{\lambda_i} |ii\rangle) = \left(\sum \sqrt{\lambda_i}\right)^2 - 1 = 2\mathcal{N}(|\phi\rangle := \sum \sqrt{\lambda_i} |ii\rangle)$$

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- ☞ **Upper bounds  $C_d$ :**

$$C_d(|\psi\rangle) = C_r(|\psi\rangle) \leq C_{\ell_1}(|\psi\rangle), \quad \forall |\psi\rangle.$$

Equality holds iff  $\text{diag}[|\psi\rangle\langle\psi|]$  is (up to permutation) either  $\{1, 0, \dots, 0\}$ , or  $\{1/2, 1/2, 0, \dots, 0\}$ . That is, either  $|\psi\rangle \in \mathcal{I}$ , or  $|\psi\rangle \equiv |\Phi_2\rangle \oplus 0$ .



Using the recursive property [Lin, IEEE Trans. Inf. Theory, 1991] of entropy function  $H(\lambda)$ , we have

$$\begin{aligned}
 C_{\ell_1} \left( |\psi\rangle = \sum_{i=1}^d \sqrt{\lambda_i} |i\rangle \right) &= C_r(|\psi\rangle) \\
 &= 2 \sum_{i=1}^{d-1} \sqrt{\lambda_i} \sum_{j=i+1}^d \sqrt{\lambda_j} - H(\lambda) \\
 &\geq 2 \sum_{i=1}^{d-1} \sqrt{\lambda_i} \sqrt{\sum_{j=i+1}^d \lambda_j} - H(\lambda) \\
 &= \sum_{i=1}^{d-1} \left[ \left( \sum_{k=i}^d \lambda_k \right) \left( 2 \sqrt{\frac{\lambda_i}{\sum_{k=i}^d \lambda_k} \left( 1 - \frac{\lambda_i}{\sum_{k=i}^d \lambda_k} \right)} - H_2 \left( \frac{\lambda_i}{\sum_{k=i}^d \lambda_k} \right) \right) \right]. \quad \square
 \end{aligned}$$

$\geq 0$ , as  $H_2(x) \leq 2\sqrt{x(1-x)}$

☞ The function  $C_{\ell_1}(|\psi\rangle) - C_r(|\psi\rangle)$  is Schur-concave in  $\text{diag}(|\psi\rangle\langle\psi|)$ .

$$0 \leq C_{\ell_1}(|\psi\rangle) - C_r(|\psi\rangle) \leq d - 1 - \log_2 d, \quad d = \text{rank}[\text{diag}(|\psi\rangle\langle\psi|)].$$

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☞ All  $|\psi\rangle$  with  $C_{\ell_1}(|\psi\rangle) = b$ , has  $C_r(|\psi\rangle)$  bounded by

$$\frac{\sqrt{2}b^2}{d(d-1)} \leq C_r \leq \log_2(1+b), \quad d = \text{rank}[\text{diag}(|\psi\rangle\langle\psi|)]$$

Note: Lower bound  $\rightarrow 0$  as  $d \rightarrow \infty$ , but the upper bound does not depend on  $d$ .

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What is optimum  $C_r$  given only the knowledge of  $C_{\ell_1}$ ?

## Optimum $C_r$ from the knowledge of $C_{\ell_1}(|\psi\rangle)$

$$H_2(\alpha) + (1 - \alpha) \log_2(d - 1) \leq C_r(|\psi\rangle) \leq H_2(\beta) + (1 - \beta) \log_2(n - 1),$$

$$\text{where } \alpha = \frac{2 + (d - 2)(d - b) + 2\sqrt{(b + 1)(d - 1)(d - 1 - b)}}{d^2},$$

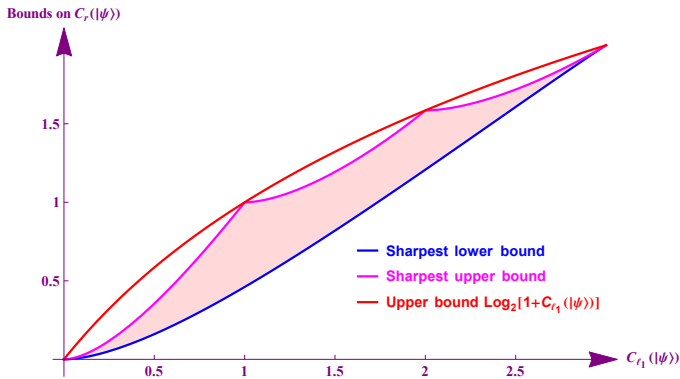
$$\beta = \frac{2 + (n - 2)(n - b) - 2\sqrt{(b + 1)(n - 1)(n - 1 - b)}}{n^2},$$

$$d = \text{rank}[\text{diag}(|\psi\rangle\langle\psi|)],$$

$$n = \begin{cases} b + 1 & \text{if } b \text{ is integer,} \\ [b] + 2 & \text{otherwise,} \end{cases}$$

with  $[x]$  denoting the integer part of  $x$ .

- Each of the bounds is satisfied by a unique state, up to permutation the diagonal elements of the state with minimum  $C_r$  are given by  $\{\alpha, (1-\alpha)/(d-1), (1-\alpha)/(d-1), \dots, (1-\alpha)/(d-1)\}$  and that with maximum  $C_r$  are  $\{\beta, (1-\beta)/(n-1), (1-\beta)/(n-1), \dots, (1-\beta)/(n-1)\}$ .
- As  $d \rightarrow \infty$ ,  $\alpha \rightarrow 1$ , so  $C_r$  could be arbitrarily small for large enough  $d$ , but can not be increased beyond that sharpest upper bound.

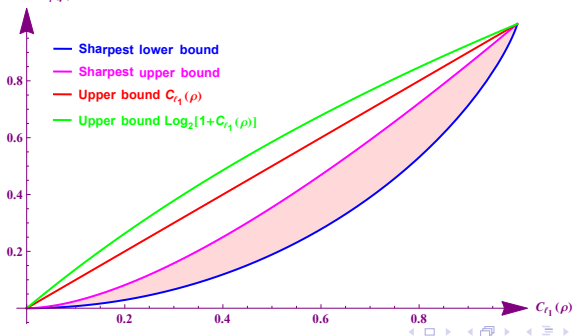


## Optimum $C_r$ from the knowledge of $C_{\ell_1}(\rho)$ for qubit $\rho$

☞ All qubit states  $\rho$  with a given coherence  $C_{\ell_1}(\rho) = 2b$  satisfy

$$1 - H_2\left(\frac{1-2b}{2}\right) \leq C_r(\rho) \leq H_2\left(\frac{1-\sqrt{1-4b^2}}{2}\right) \leq C_{\ell_1}(\rho).$$

Bounds on  $C_r(\rho)$



## Optimum $C_r$ from the knowledge of $C_{\ell_1}(\rho)$ for general $\rho$ ?

All we have is:

$$C_r(\rho) \leq \log_2 [1 + C_R(\rho)] \leq \log_2 [1 + C_{\ell_1}(\rho)] \leq \begin{cases} C_{\ell_1}(\rho), & \text{if } C_{\ell_1}(\rho) \geq 1 \\ \frac{1}{\ln 2} C_{\ell_1}(\rho), & \text{if } C_{\ell_1}(\rho) < 1 \end{cases},$$

where  $C_R$  is the robustness, a strong monotone, defined by

$$C_R(\rho) := \min_{\sigma} \left\{ s \geq 0 \mid \frac{\rho + s\sigma}{1+s} \in \mathcal{F} \right\} = \min_{\tau \in \mathcal{F}} \{ s \geq 0 \mid \rho \leq (1+s)\tau \}.$$



## Conclusion and Outlook

- Both  $\ell_p$ - and Schatten- $p$ -norm based functions are not strong monotone of coherence for any  $p > 1$ . Neither is a monotone in the first place.  $C_{\text{Tr}}$  is monotone, but not strong monotone like  $C_{\ell_1}$ .

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- $E_{\text{Tr}}$  is also monotone but not strong monotone. In [Eisert *et al.* JPA 2003] it is claimed that if we restrict the minimization over separable states having the same reductions as  $\rho$ , then the modified  $E_{\text{Tr}}$  is a strong monotone. What is an easy modification for  $C_{\text{Tr}}$  to make it a strong monotone?

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- It seems that like convexity (as logarithmic version of any strong monotone is non-convex), strong monotonicity should be only a desirable feature for monotone, not a primary requirement.

- Operational (information theoretic) interpretations of  $C_{l_1}$ :

$$C_d(|\psi\rangle) = C_r(|\psi\rangle) \leq C_{l_1}(|\psi\rangle), \forall |\psi\rangle$$

$$C_d(\rho) = C_r(\rho) \leq \log_2 [1 + C_{l_1}(\rho)] \leq \begin{cases} C_{l_1}(\rho), & \text{if } C_{l_1}(\rho) \geq 1 \\ \frac{1}{\ln 2} C_{l_1}(\rho), & \text{if } C_{l_1}(\rho) < 1 \end{cases} .$$

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Question: Can we drop the factor  $1/\ln 2 \approx 1.4427$ ?

- If so, it would be sharpest as for any  $0 < b < 1$  and  $d \geq 3$ , the  $d$ -dimensional state

$$\rho = \begin{pmatrix} b/2 & b/2 \\ b/2 & b/2 \end{pmatrix} \oplus (1-b)\delta, \quad \delta \in \mathcal{F}$$

has  $C_r(\rho) = C_{\ell_1}(\rho) = b$

# Thank You!